

Detectability of cold streams in absorption and emission



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Cold streams

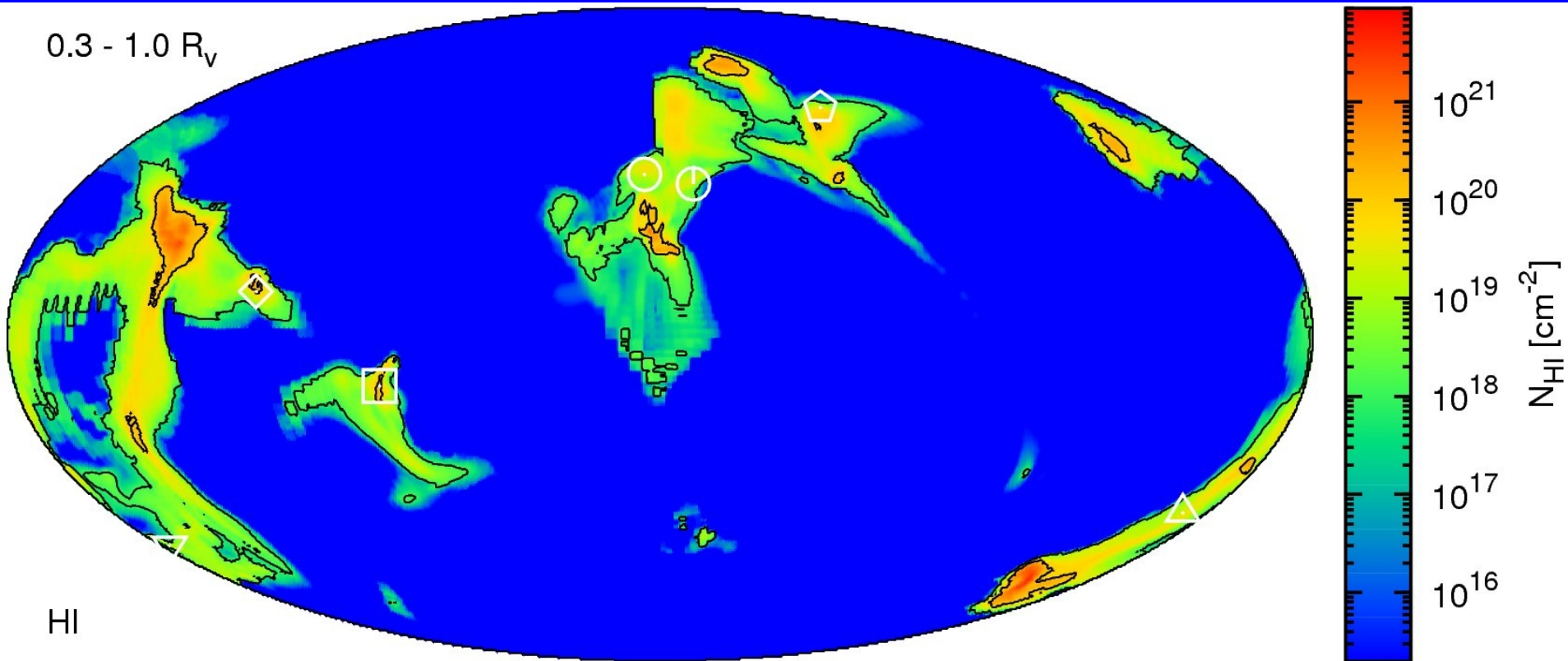


Steidel et al. 2010:

- Observes “Circum Galactic Medium”
- Absorption line profiles
- Stacks more than 100 spectra
- Detects massive outflows
- But: no sign of inflows
- Claim: Proof of absence of cold streams

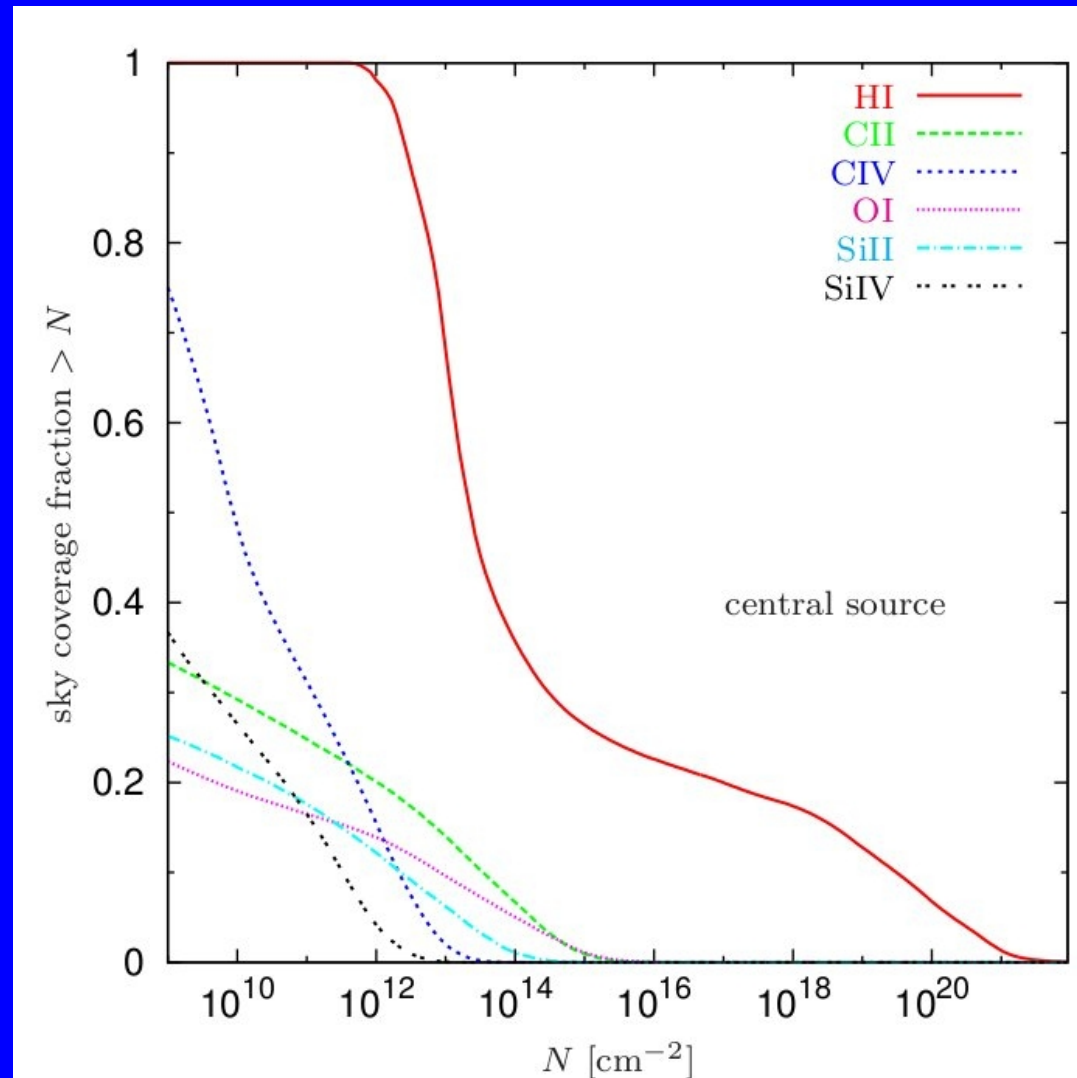
Central geometry

- Observes central galaxy through its own circum galactic medium
- AMR simulation, $z = 2.5$, $M_{\text{vir}} = 4 * 10^{11} M_{\odot}$



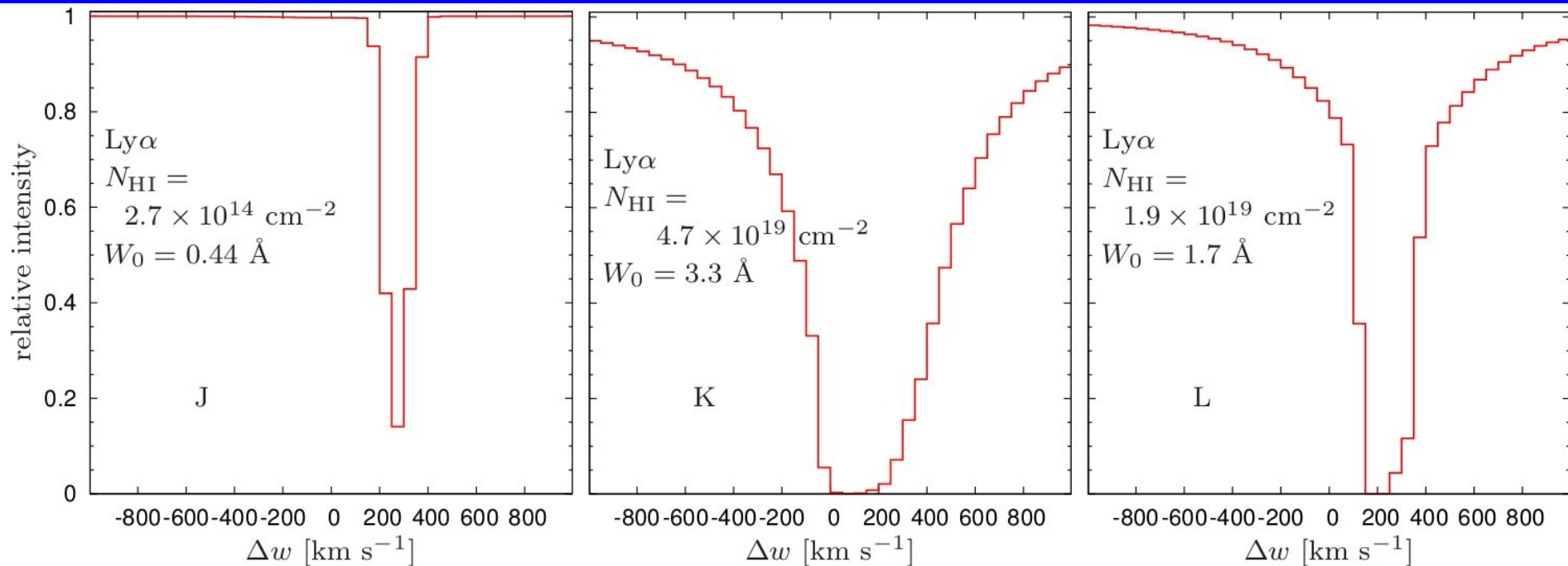
Sky covering fraction

- Very low sky covering fraction
- Low metallicity in streams



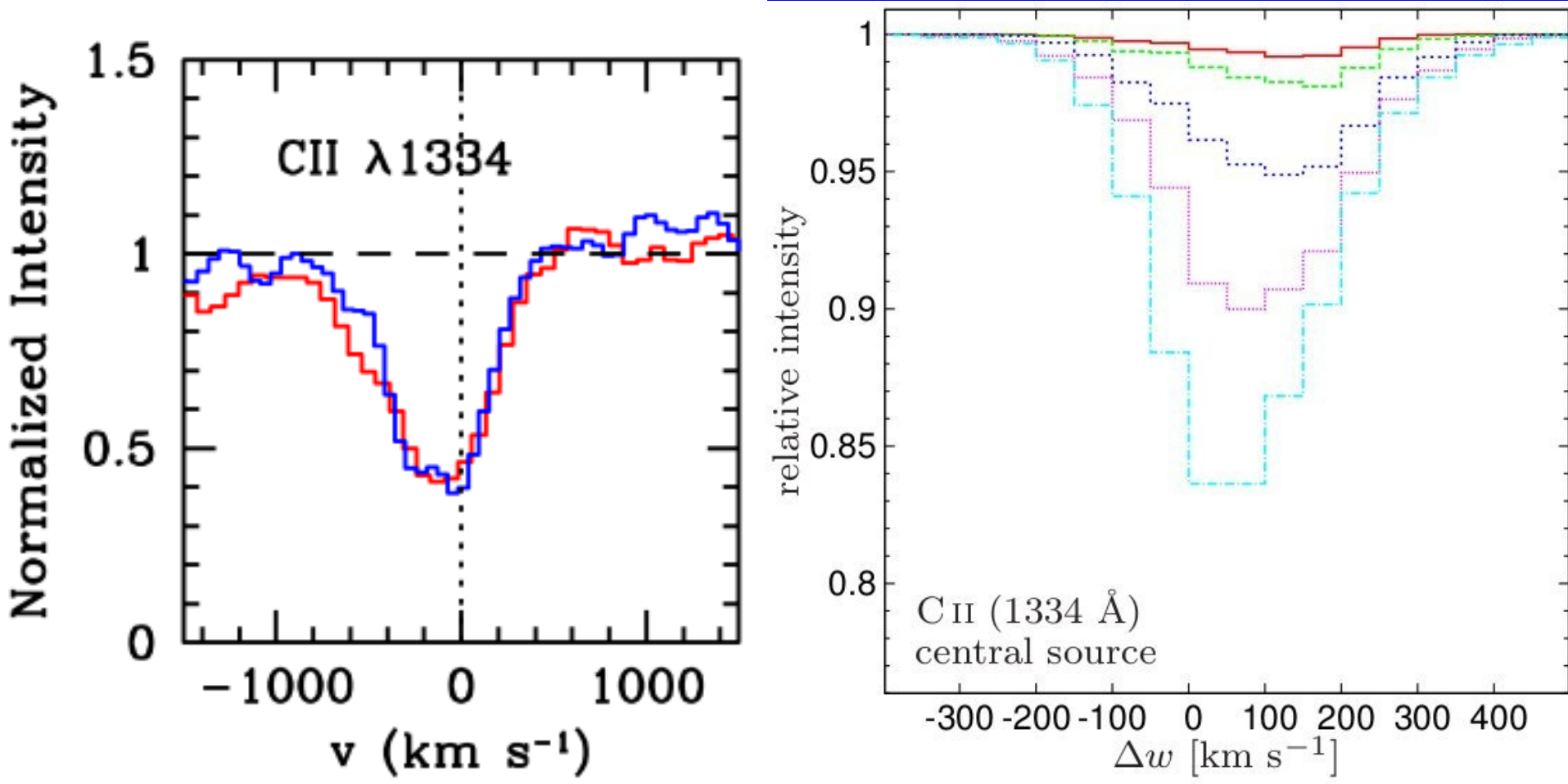
Example line profiles

- Ly α
- Gaussian point spread function with 4kpc beam-size applied
- Velocity resolution degraded to 50 km s $^{-1}$
- Observer convention: inflow positive (right)



Stacked line profile:

- Averaging over all available example line profiles (3 galaxies, all directions)

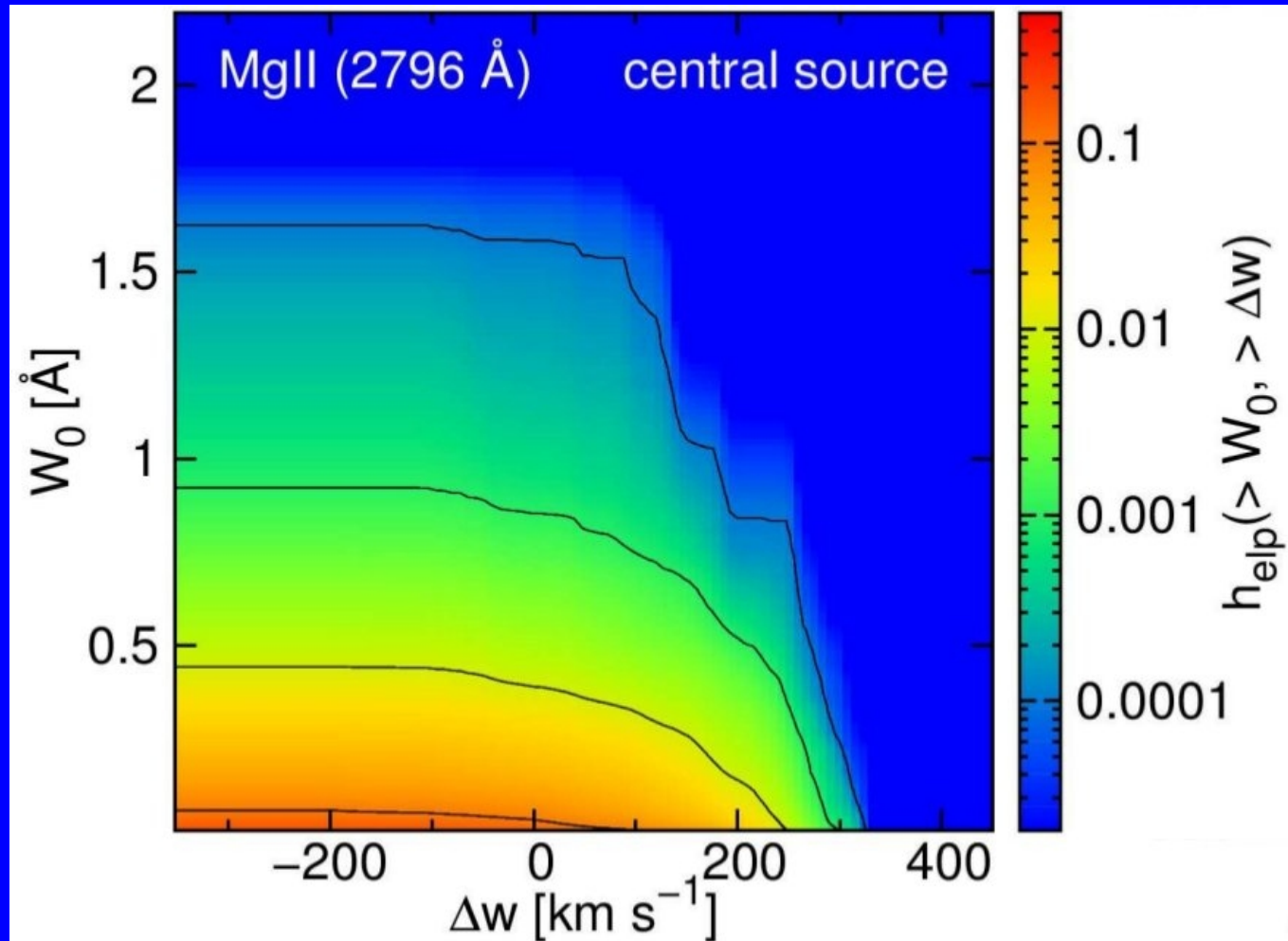


Stacking?

- Stacking washes out the cold filament absorption signal
- Cold filament absorption signal might still be visible in non stacked data

Statistics

- Mg II: inflow $> 150 \text{ km s}^{-1}$ with an EW $> 0.2 \text{ \AA}$ in 1.3 % of all observations

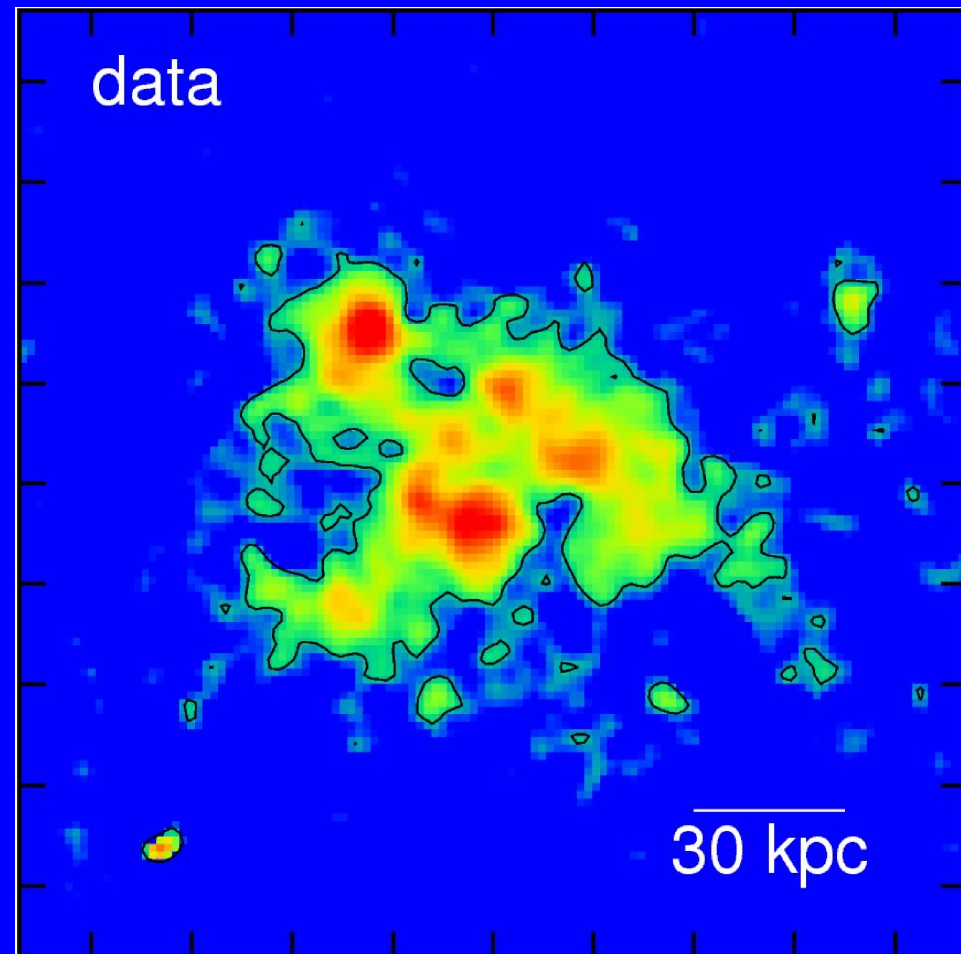


Absorption summary

- Observational features from cold streams extremely difficult to detect.
- Outflows are dominant.
- No falsification done.

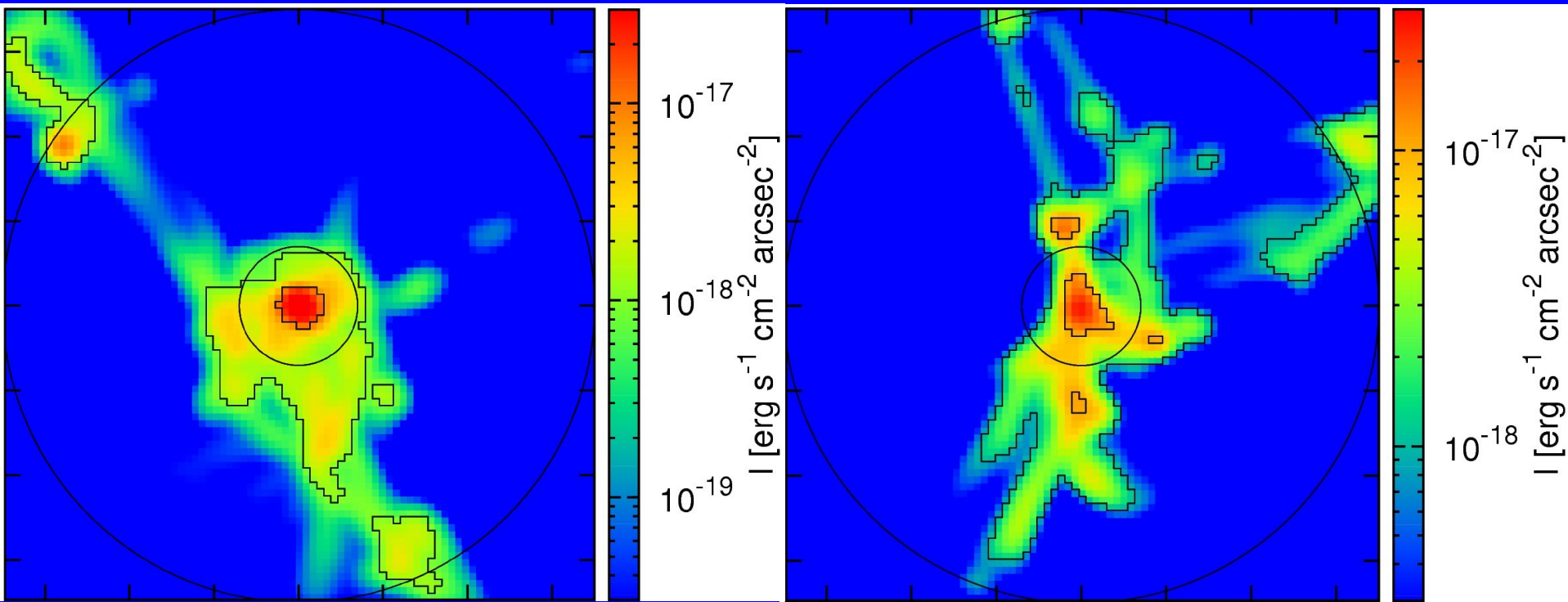
Lyman alpha blobs

- First observed by Steidel et al. 2000
- Redshift range $z = 2 - 6.5$
- Observation by Matsuda et al. 2004

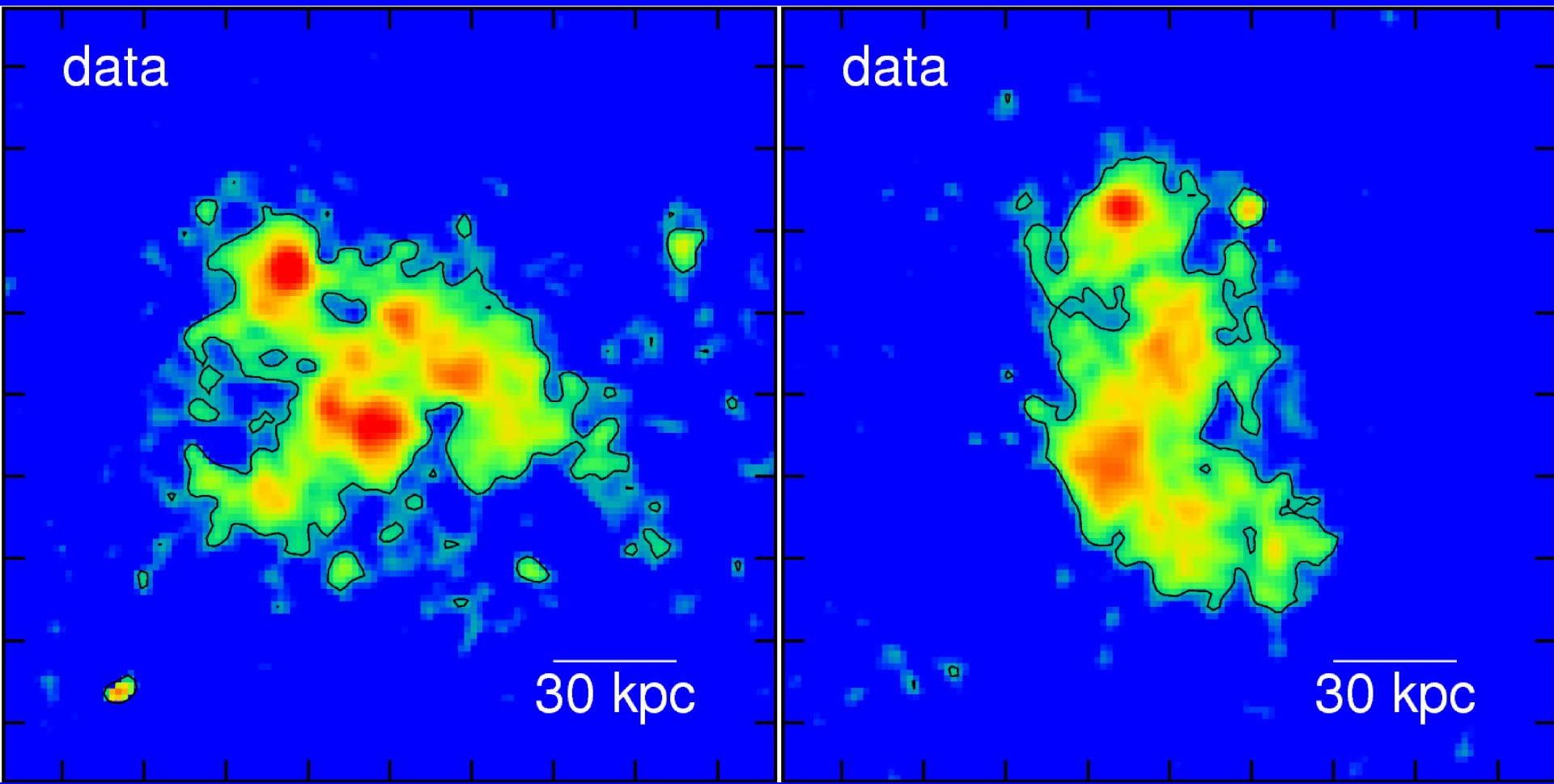


Simulated surface brightness maps:

- AMR simulation (CDB)
- $Z = 3.09$
- $M_{\text{vir}} = 3.5e11 M_{\odot}$
- 0.6" FWHM Gaussian PSF

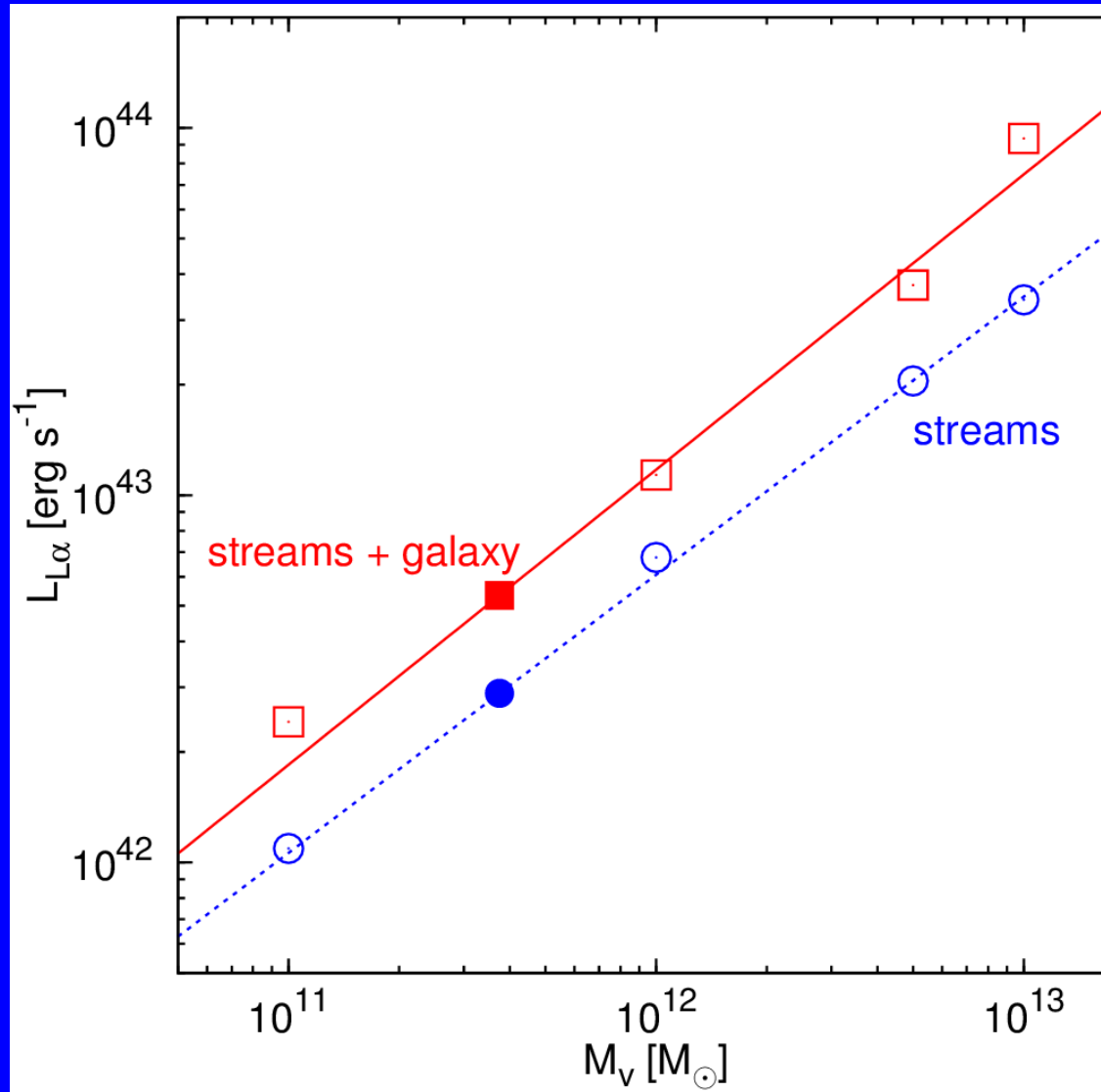


The observations for comparison:



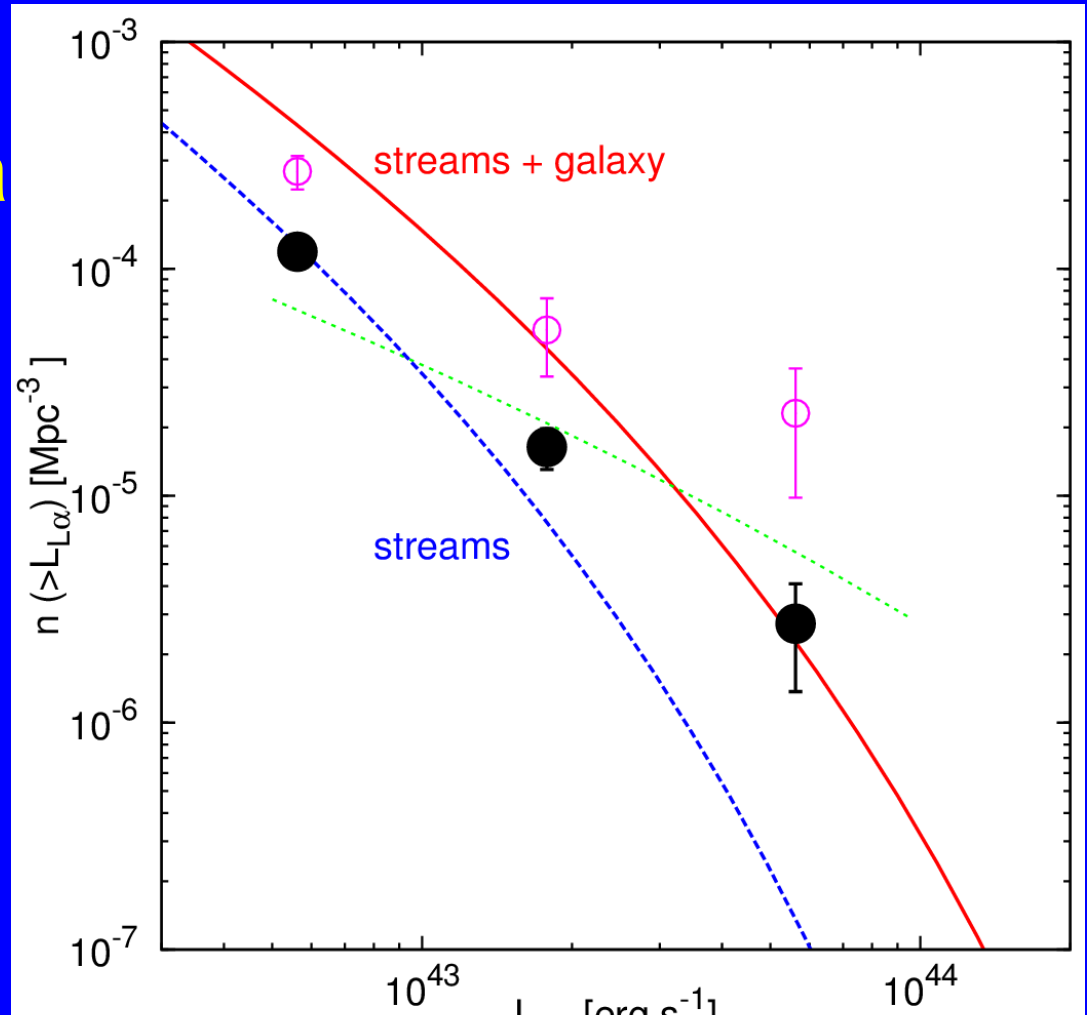
Lyman alpha vs halo mass

- AMR simulations
- Several galaxies per data point
- $z = 3.09$



Luminosity function

- Correlation with Sheth Tormen mass function
- Data from Matsuda et al. 2004



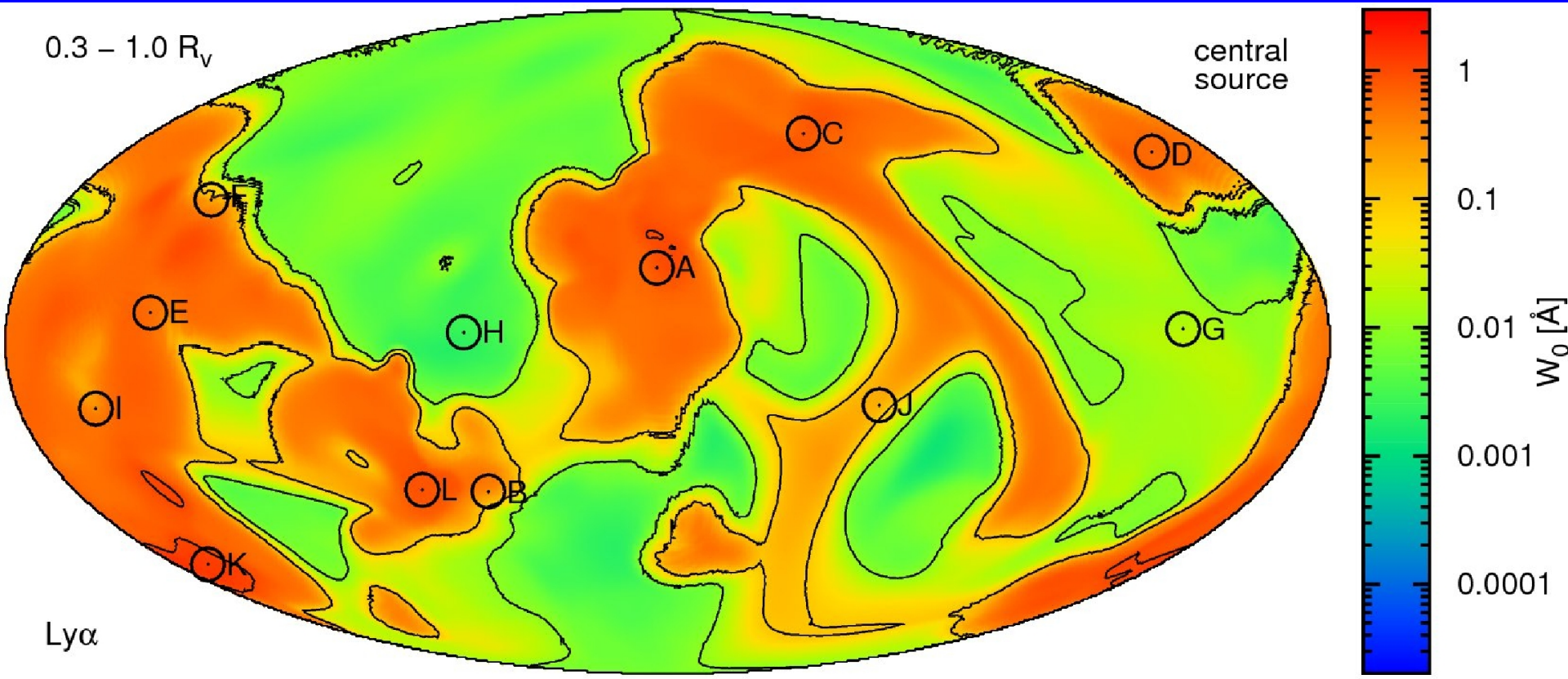
Emission summary

- Cold streams loose pot. energy released as Lyman alpha photons
- Simulation maps very similar to observations in extent, shape, luminosity
- Luminosity function fits data
 - => Cold streams can explain Lyman alpha blobs
 - => First observational evidence for cold streams!

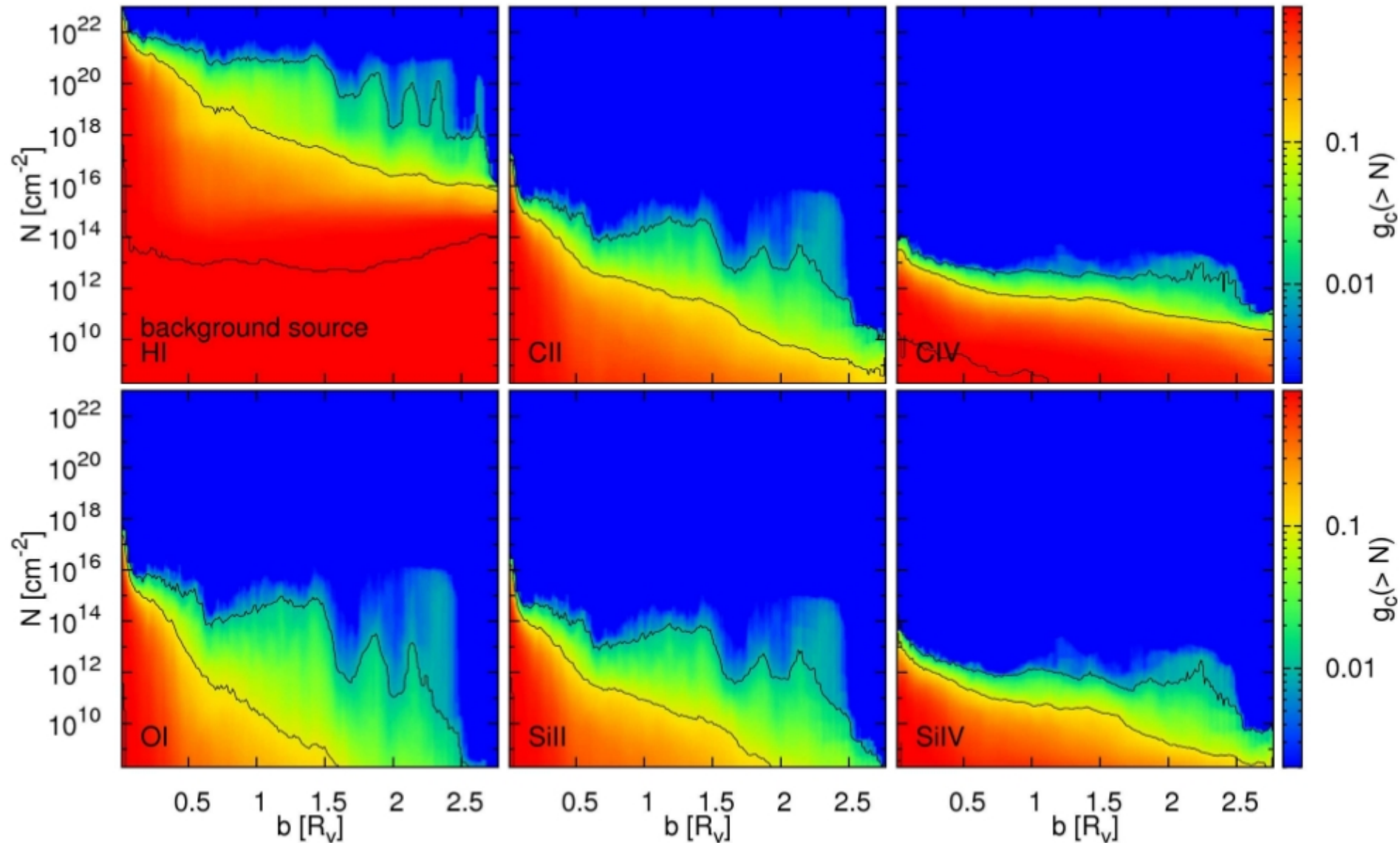
Thanks!

Equivalent width

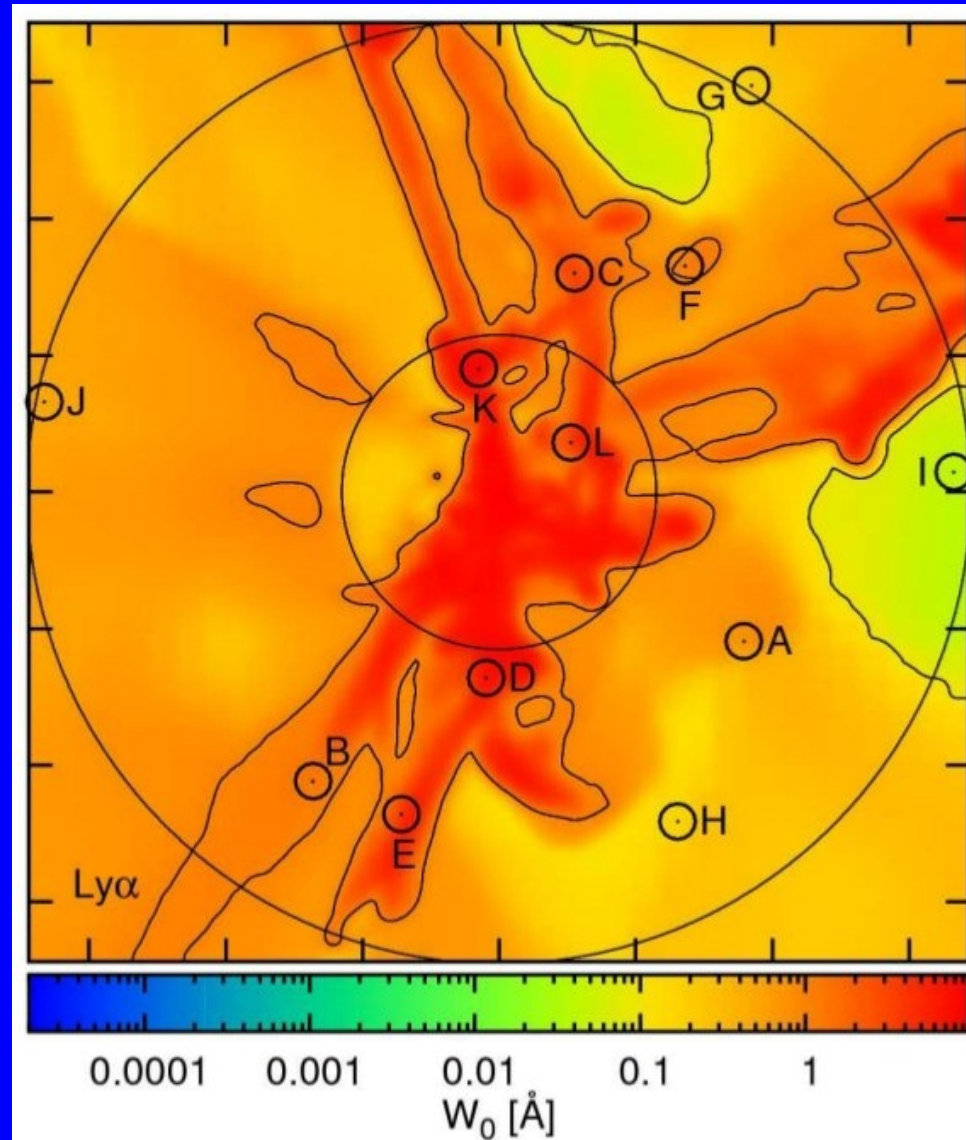
$$EW = \frac{\lambda^2}{c} \int_0^\infty [1 - \exp(-\tau_\nu)] d\nu = \frac{\lambda^2}{c} \int_0^\infty g(\nu) d\nu,$$



covering fraction vs impact parameter vs column density

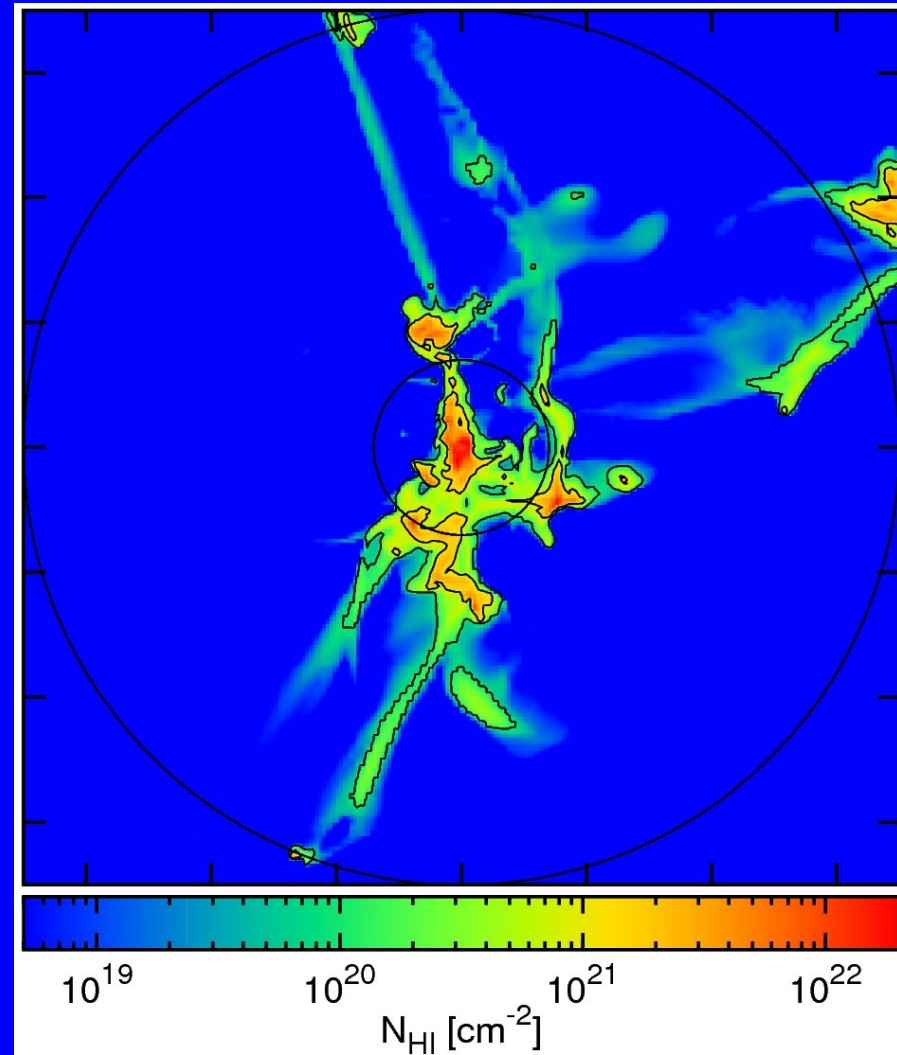


Equivalent width



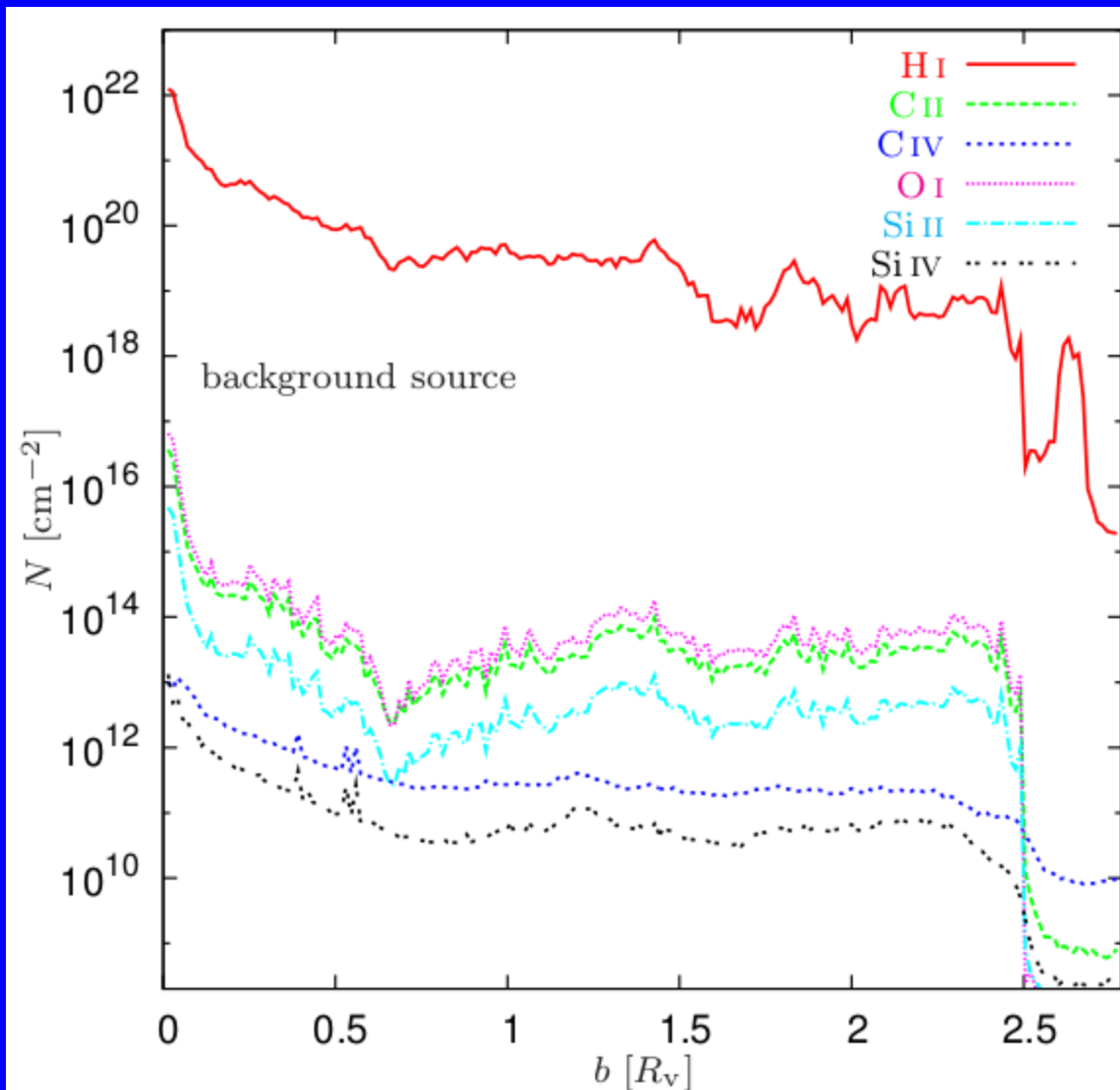
Background geometry

- Observes background galaxy through circum galactic medium of galaxy in question
- Additional parameter: Impact parameter b



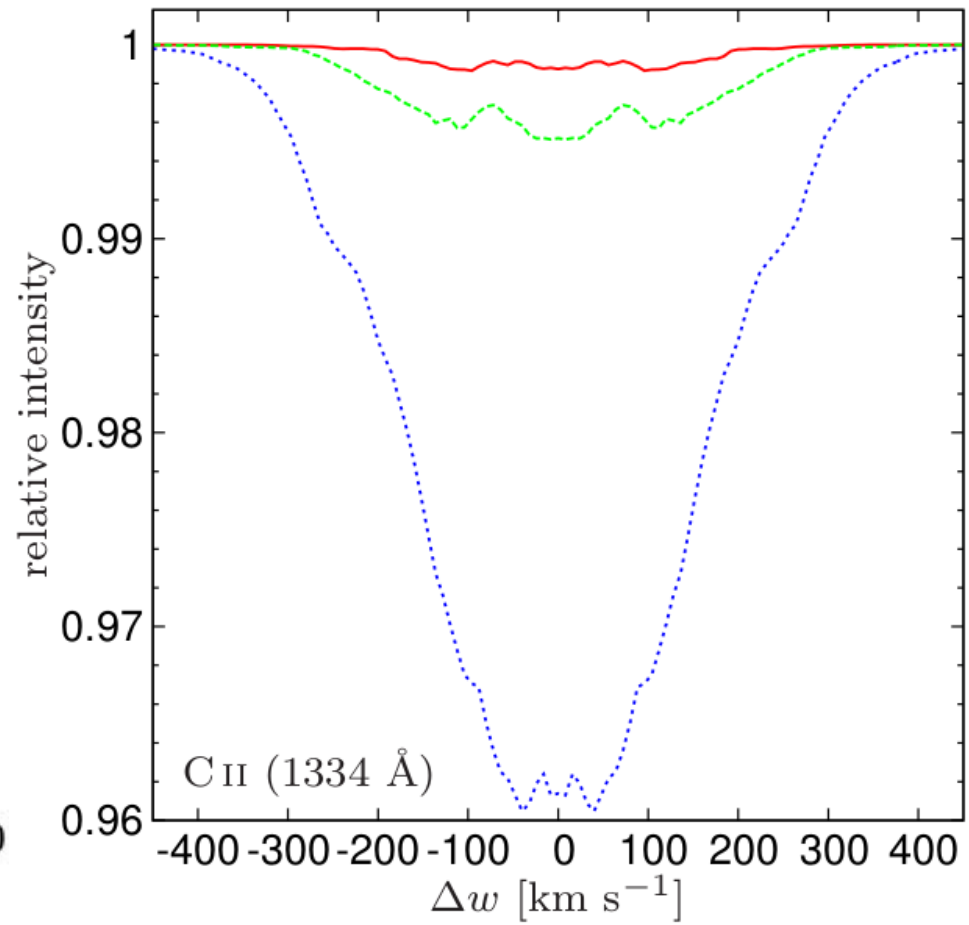
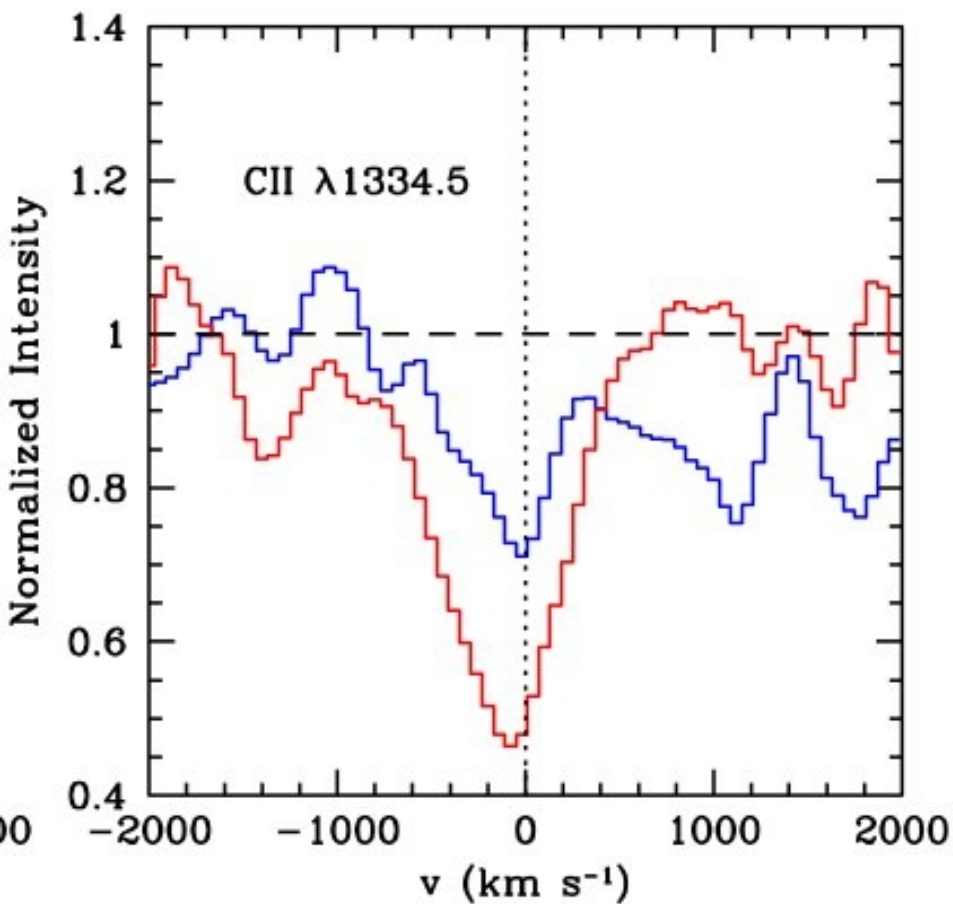
Column density vs impact parameter

- All lines decreasing
- Ly alpha considerably higher than metals



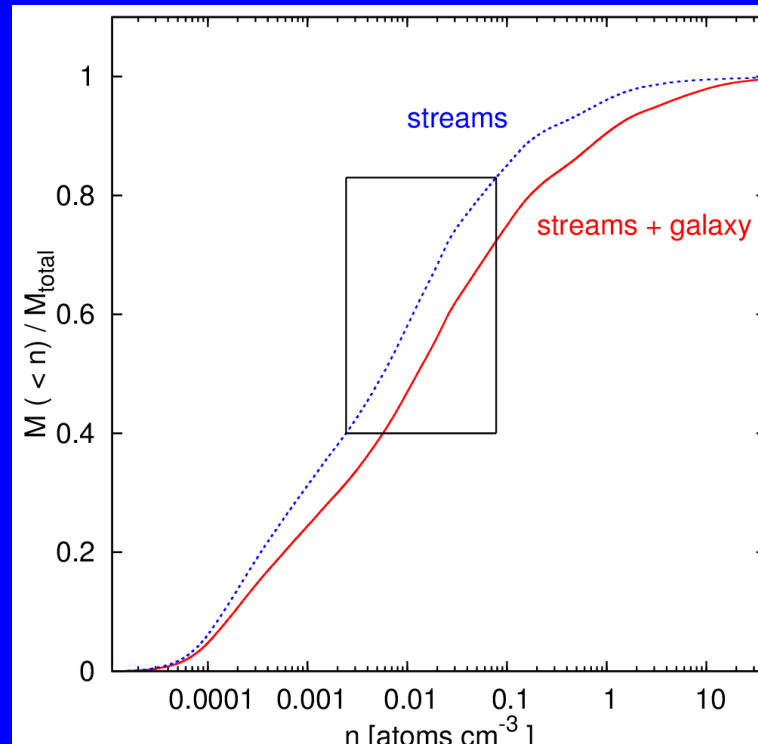
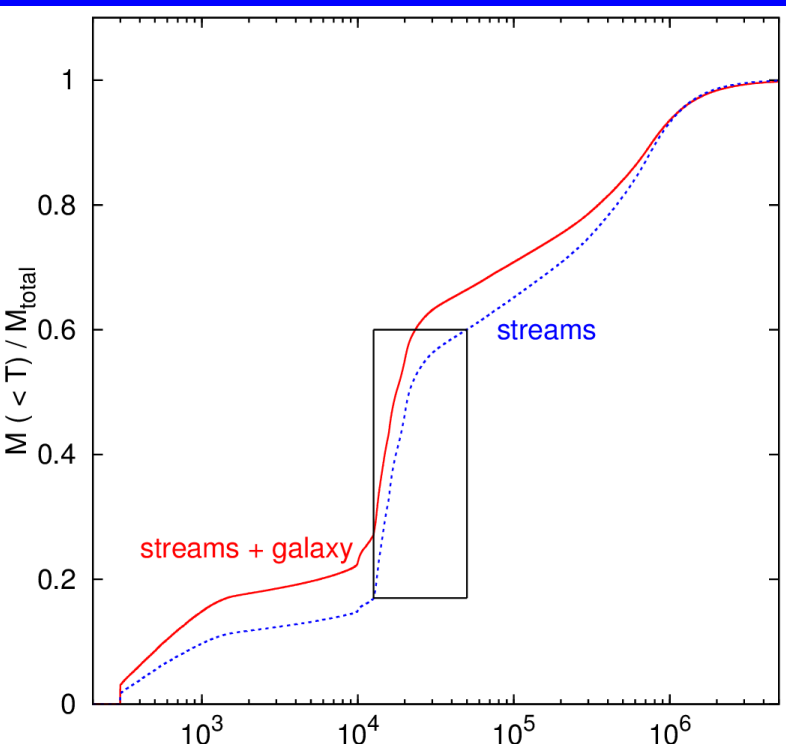
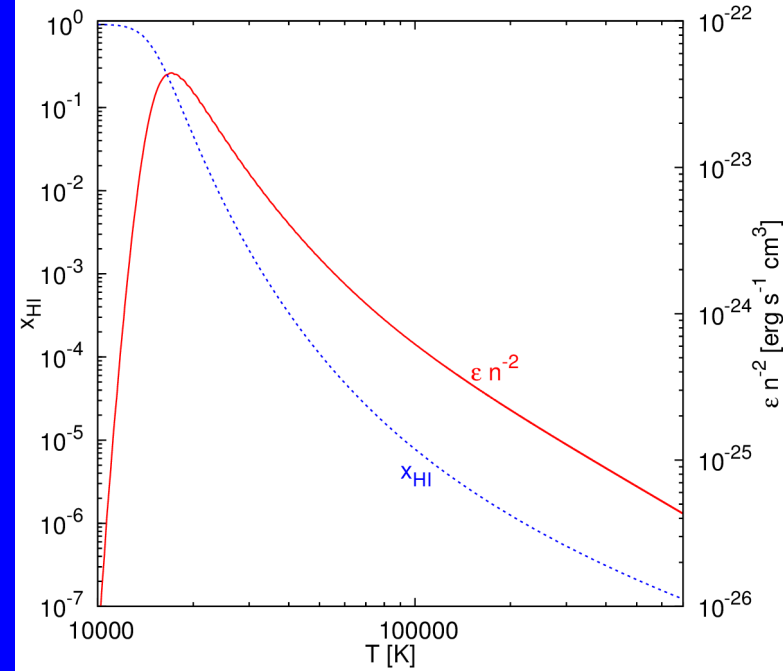
Stacked line profile

- Averaging over all available example line profiles (3 galaxies, 6 principal directions, all points in radiusrange)



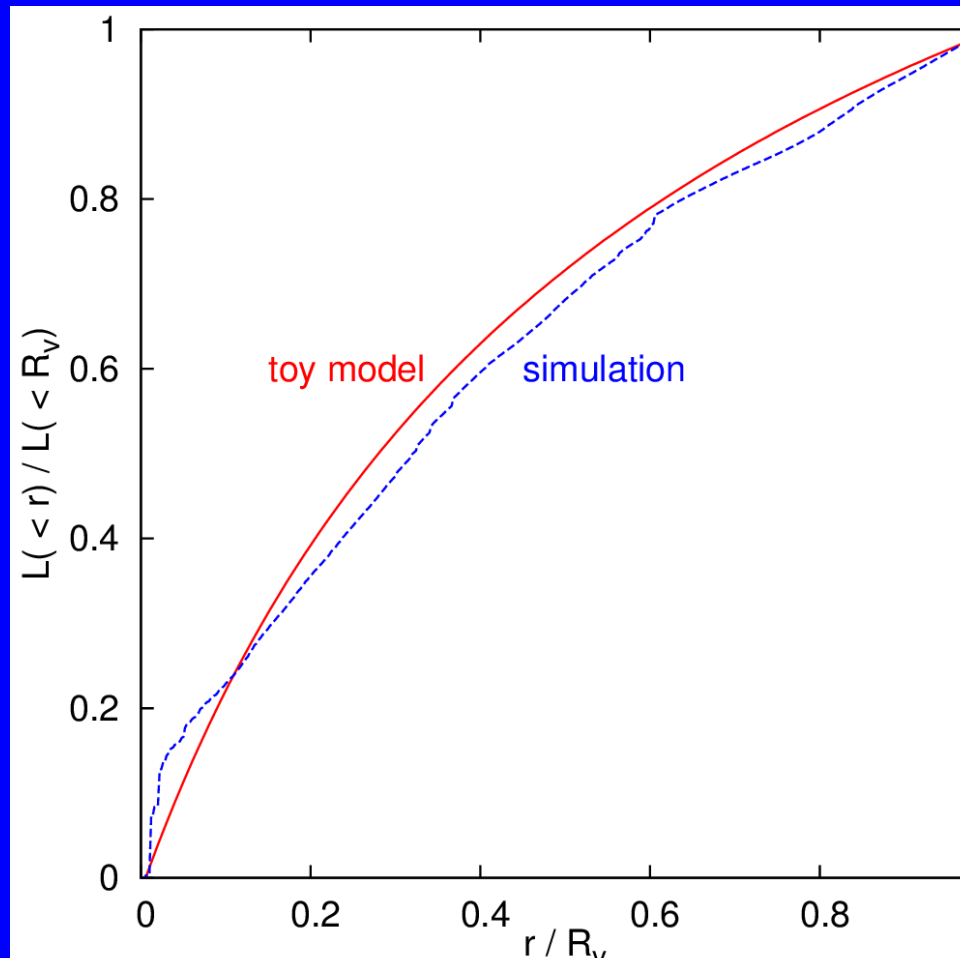
L α emissivity:

- 50% of the gas emits L α efficiently



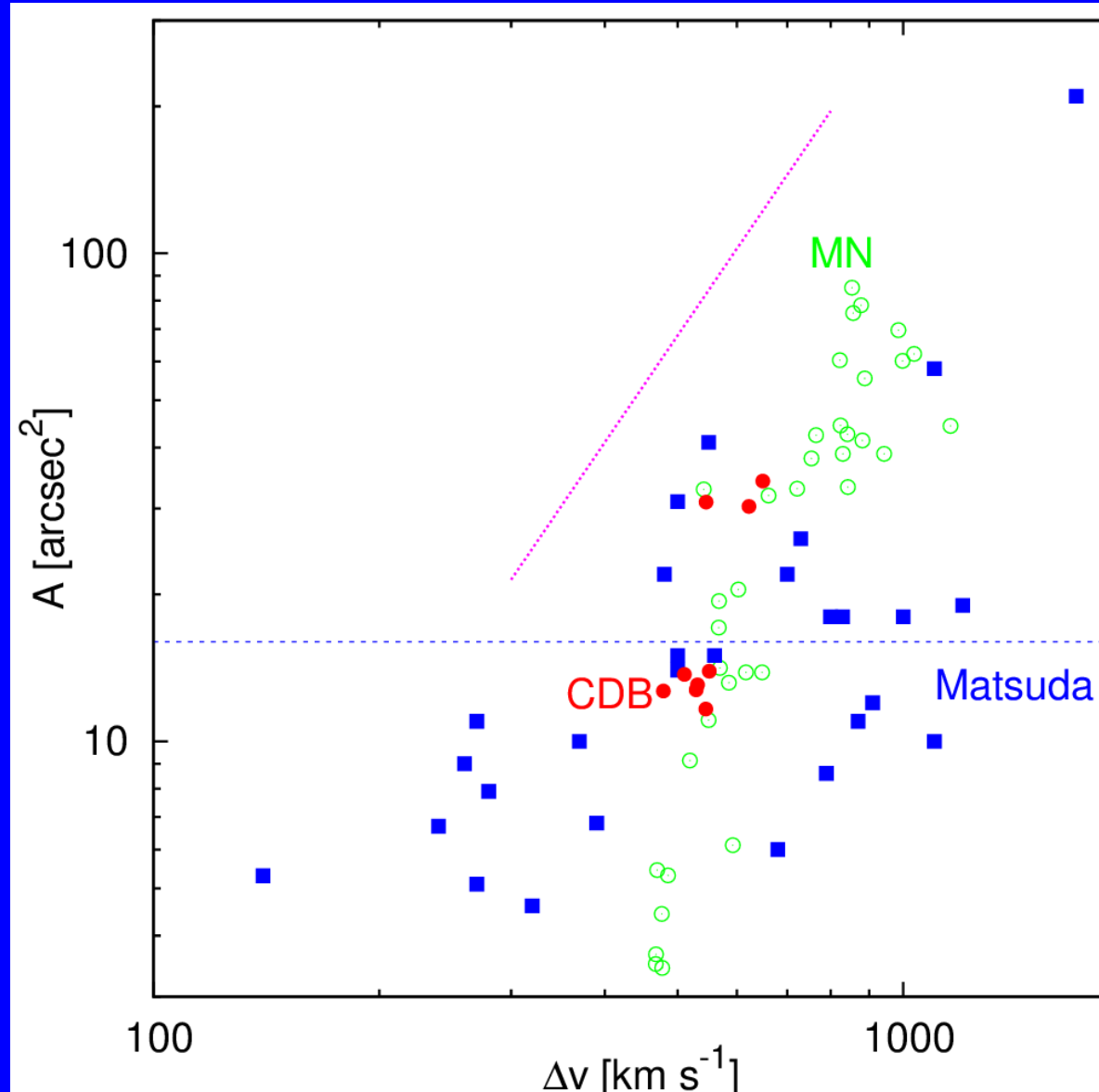
Toy model:

- NFW profile
- Neistein infall (EPS)
- Constant infall velocity



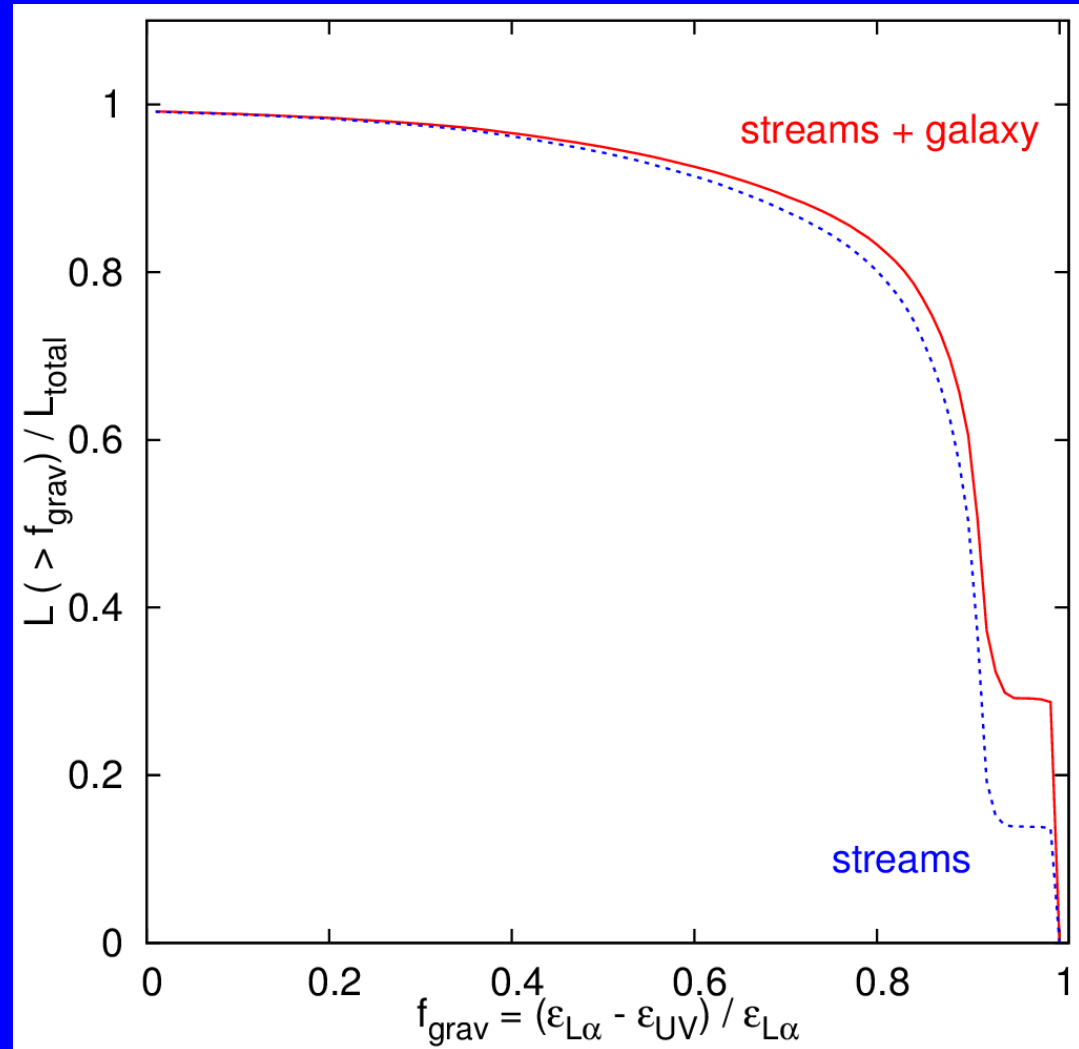
Kinematics

- Area vs. velocity dispersion



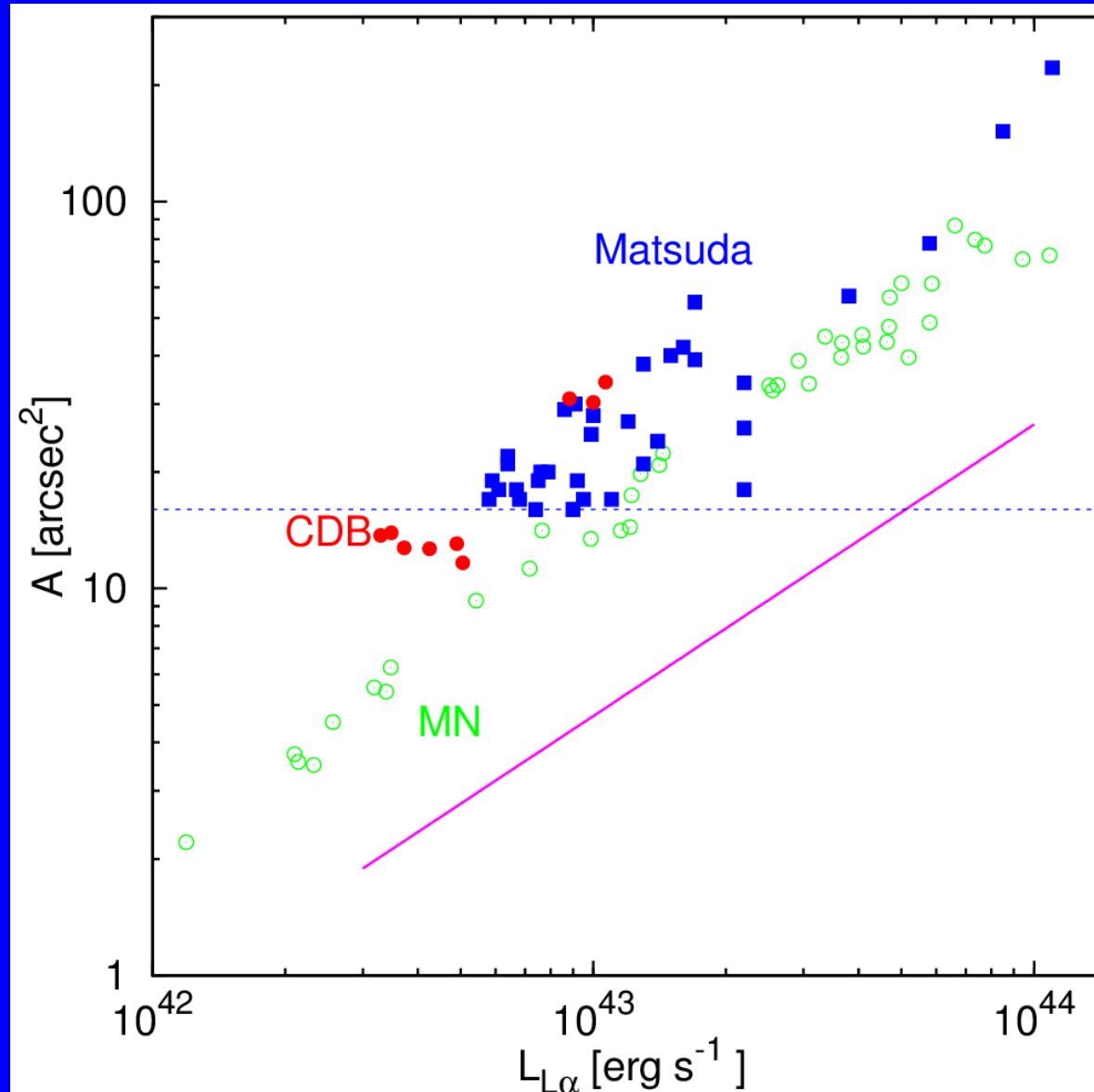
Energy source: Gravitational heating vs. UV background

In the gas that contributes 80% of the luminosity more than 80% of the input energy is gravitational



Area vs. Luminosity

- Isophotal area above 2.2×10^{-18} $\text{erg s}^{-1} \text{cm}^{-2}$ arcsec^{-2} as a function of total luminosity



The AMR simulations

- Ceverino, Dekel & Bournaud
 - Art by Andrey Kravtsov
 - UV background, Haardt & Madau 1996
 - mimics self-shielding
 - Gas can cool down to 100K
 - 3 re-simulated galaxies
 - High resolution (70 pc physical)

Computing Lyman alpha

- Emissivity:

$$\epsilon = n_e n_{\text{HI}} C_{\text{L}\alpha}(T) + 0.68 h\nu_{\alpha} n_e n_{\text{HII}} \alpha_{\text{rec,B}}(T)$$

- Collisional excitation coefficient:

$$C_{\text{L}\alpha} = 3.7 \times 10^{-17} T^{-1/2} \exp\left(-\frac{h\nu_{\alpha}}{kT}\right) \text{ erg s}^{-1} \text{ cm}^3$$

- Case-B recombination coefficient:

$$\alpha_{\text{rec,B}}(T) = 4.9 \times 10^{-6} T^{-1.5} \left(1 + \frac{115}{T^{0.41}}\right)^{-2.24} \text{ cm}^3 \text{ s}^{-1}$$

More computing

- Number densities:

$$n_{\text{HI}} = \frac{x_{\text{HI}} X \rho}{m_{\text{p}}},$$

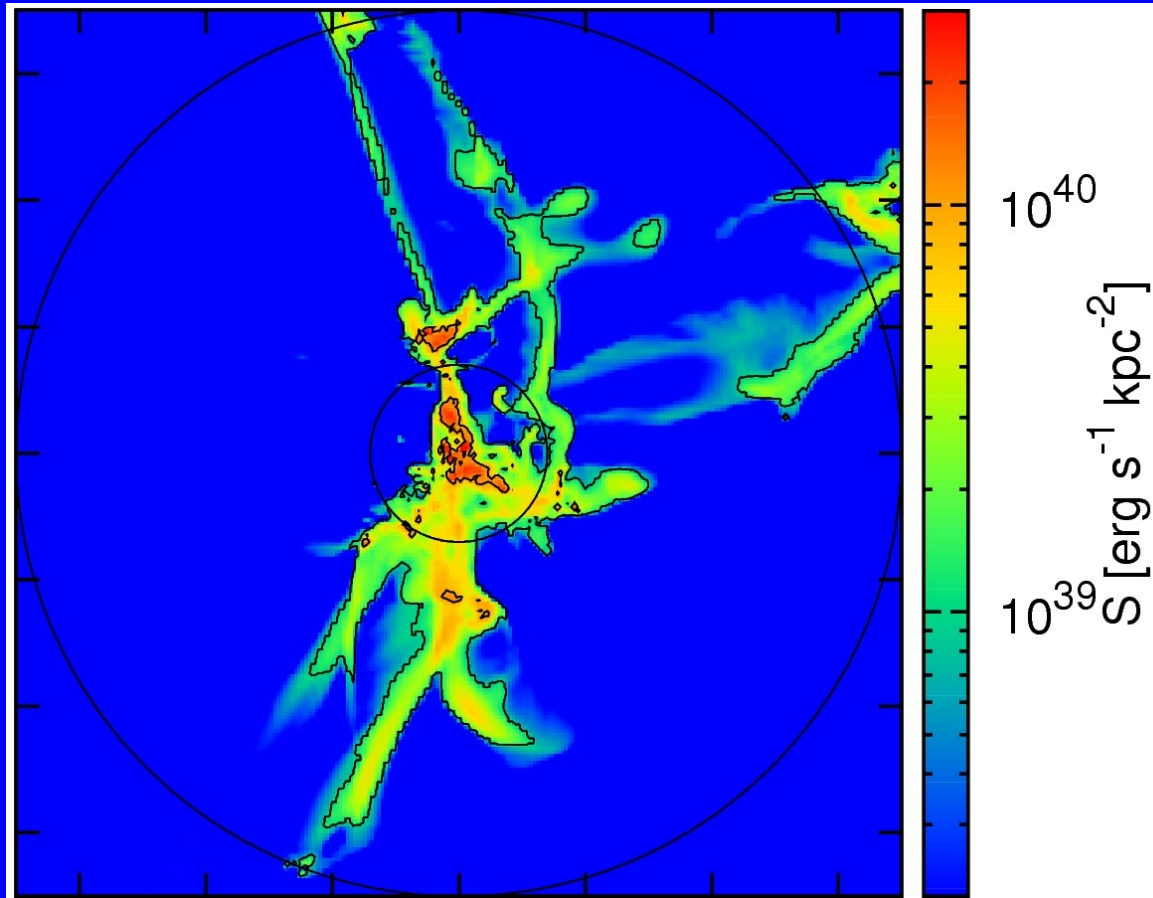
$$n_{\text{HII}} = n_e = \frac{(1 - x_{\text{HI}}) X \rho}{m_{\text{p}}}$$

- Neutral Hydrogen fraction:

$$x_{\text{HI}} = \frac{\alpha_{\text{rec,B}}(T)}{\alpha_{\text{rec,B}}(T) + C_{\text{ion}}}$$

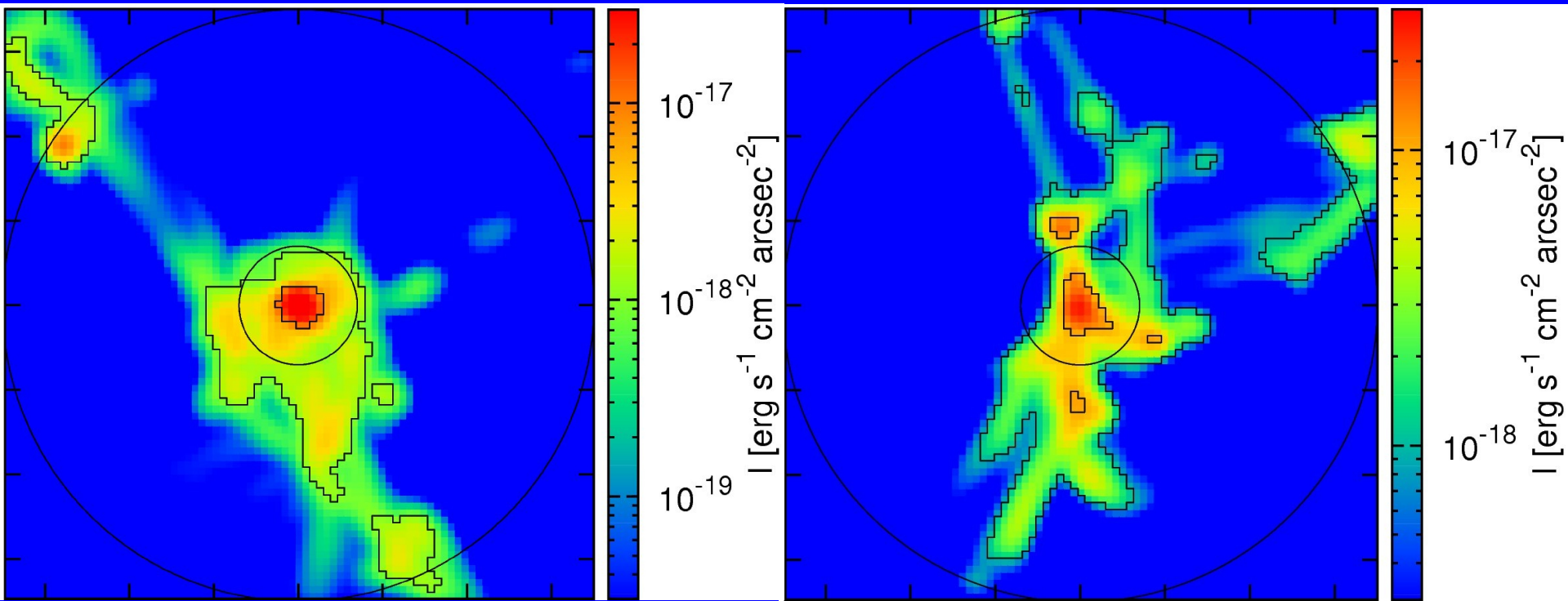
Resulting maps: Surface brightness

- CDB simulation
- $Z = 3.09$
- $M_{\text{vir}} = 3.5e11 M_{\odot}$



What an observer would see:

- $I = S / [4 \pi (1+z)^4]$
- 0.6" FWHM
Gaussian PSF



Computing line profiles

- Doppler broadening

$$b = \sqrt{\frac{2kT}{m_Y}},$$

- Optical depth τ

$$\begin{aligned} \tau_\nu(\phi, \theta, \Delta\omega) &= \frac{\sqrt{\pi} e^2 f_\lambda \lambda_0}{m_e c} \int_{r_i}^{R_\nu} \frac{n_Y(\vec{r}) X_{XX}(\vec{r})}{b(\vec{r})} \\ &\times H \left[\frac{\gamma_\lambda \lambda_0}{4\pi b(\vec{r})}, \frac{\Delta\omega - v(\vec{r})}{b(\vec{r})} \right] dr, \end{aligned}$$

- Intensity $I(\Delta\omega) = \exp(-\tau)$

Additional AMR simulations

- Horizon MareNostrum
 - Ramses by Romain Teyssier
 - UV background: Haardt & Madau 1996
 - Density dependent pressure floor:
 - $T_{\text{floor}} = 10^4 (n/0.1)^{2/3} \text{K}$ for $n > 0.1 \text{cm}^{-3}$
 - Fully cosmological simulation
 - Fairly good resolution (1kpc physical)